Lessons From Last Week #1

Interfaces of Two-Type Continuum-Sites Stepping-Stone Models

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The infinite-density, ordered model embeds finite-density, symmetric models of all densities.

If

$$X_t^K = \sum_{u \in \mathfrak{U}, \, u < K} \delta_{\left(\widehat{Z}_u(t), \kappa_u(t)\right)}$$

 $\Psi_t = \sum_{u \in \mathfrak{U}} \delta_{\left(\hat{Z}_u(t), \kappa_u(t), u\right)}$

then

 $X^K = {}^d K$ -density symmetric model

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Lessons From Last Week #2

The limiting location type measure

$$X_t = \lim_{K \to \infty} \frac{1}{K} X_t^K = \lim_{K \to \infty} \frac{1}{K} \sum_{u \in \mathfrak{U}, u < K} \delta_{(\hat{Z}_u(t), \kappa_u(t))}$$

- "exists" as an L_p and a.s. limit;
- is the weak limit of the (mass-scaled)
 K-density, symmetric model as K → ∞;
- satisfies

$$\mathfrak{L}\left(\Psi_t \mid \mathfrak{F}_t^X\right) = \mathsf{Poisson}(X_t \times \ell_{\mathbb{R}_+})$$

The "essence" of the infinite-density, continuum-sites, stepping-stone model is X_t . It is generated by (or generates) a conditionally Poisson particle system. Two-Type, No Mutation Case

Let $E = \{0 \equiv \bigcirc, 1 \equiv \bullet\}$ with no mutation.

Then $X_t(\cdot \times \{1\}) \leq X_t(\cdot \times E) = \ell_{\mathbb{R}}$, so let $X_t(dx \times \{1\}) = u_t(x)dx$.



$$\dot{u} = \frac{1}{2}\Delta u + |2\gamma u(1-u)|^{1/2} \dot{W}$$

That is, writing $u_t(f) = \int f(x)u_t(x)dx$, we have

$$u_t(f) = u_0(f) + \frac{1}{2} \int_0^t u_t(f'') ds + M_t^f$$
$$[M^f]_t = 2\gamma \int_0^t \int_{\mathbb{R}} f^2(x) u_s(x) \left(1 - u_s(x)\right) dx ds$$

This is the *X*-measure for the infinite-density model with

- two types, no mutation;
- ind BMs at rate $\theta = 1$;
- ordered pair interacting at $\frac{\gamma}{2}dL_t^0(\hat{Z}_u-\hat{Z}_v)$.

Space-Time Scaling (Tribe 1995)

$$\dot{u} = \frac{1}{2}\Delta u + |2\gamma u(1-u)|^{1/2} \dot{W}$$

Start all zero to left of origin, all ones to right: $u_0 = 1_{[0,\infty)}$.

Define space-time scaling:

$$v_t^{(n)}(x) = u_{n^2t}(nx)$$

In the limit $n \to \infty$, there is a one-point interface whose motion converges to rate 1 Brownian motion.

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Another View of the Scaling



$$\dot{u} = \frac{1}{2}\Delta u + |2\gamma u(1-u)|^{1/2} \dot{W}$$

then $v_t^{(n)}(x) = u_{n^2t}(nx)$ solves

$$\dot{v}^{(n)} = \frac{1}{2}\Delta v^{(n)} + \left|2n\gamma v^{(n)}\left(1-v^{(n)}\right)\right|^{1/2}\dot{W}$$



- Does $\gamma\text{-lookdown system converge to}$ immediate-lookdown system as $\gamma \to \infty?$
- Is $g \circ f$ a "continuous" function of the particle system?

Immediate-Interaction Model



- at time 0, all to left, all to right of origin;
- independent, rate 1 horizontal Brownian motions;
- immediate lookdowns: $\hat{Z}_u(t) = \hat{Z}_v(t)$ implies $\kappa_u(t) = \kappa_v(t)$.

Reference: Donnelly, Evans, Fleischmann, Kurtz, and Zhou. "Continuum-sites stepping-stone models, coalescing exchangeable partitions, and random trees." *Annals of Probability*, 28(3):1063–1110, 2000.

Interface Persists

We never get a \bigcirc to the right of a \bigcirc because of the...

Ancestral Trajectory of Particles

Track a particle at time t = T backward to t = 0. Follow each lookdown.



Note that the type can't change: it can only change at a lookdown, and we follow those.

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Interface Persists

Suppose at time t = T, there's a \bigcirc to the right of a \bigcirc . Then:

- follow ancestral trajectories back to t = 0;
- at t = 0, \bigcirc is left of \bigcirc , so they crossed;
- at crossing, there's a lookdown making them the same type, a contradiction.

Let's write:

$$L_t = \sup\{\hat{Z}_u(t) : \kappa_u(t) = 0 \equiv 0\}$$

$$R_t = \inf\{\hat{Z}_u(t) : \kappa_u(t) = 1 \equiv \bullet\}$$

and note $L_t \leq R_t$.

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Interface Is a Point

At each time *t*, particles are $Poisson(\ell_{\mathbb{R}} \times \ell_{\mathbb{R}_+})$ in location/level space.

Suppose interface is larger than a point:



There are an infinite number of particles between L_t and R_t . So what are their types?

Lemma 1. For all a < b and $t \ge 0$ a.s., there are an infinite number of particles in (a, b) in a neighborhood of t.

Interface Process Is Continuous

Lemma 2. For all a < b and $t \ge 0$ a.s., in a neighborhood of t, a particle confined to (a,b) does not change its type.

Fix $t \ge 0$. For any $\epsilon > 0$, take intervals $(L_t - \epsilon, L_t)$ and $(L_t, L_t + \epsilon)$.



In *t*-nbd given by lemma, $|L_r - L_t| < \epsilon$.

Interface Inherits Structure

Lemma 3. For all $t \ge 0$,

 $\mathfrak{L}\left(\Psi_{t} \mid \mathfrak{F}_{t}^{L}\right) = \Pi_{L_{t}}$

where $\Pi_x \equiv \text{Poisson} \left(\ell_{(-\infty,x)} \times \delta_0 + \ell_{[x,\infty)} \times \delta_1 \right).$



Proof of Lemma 3

Enough to show that $\mathfrak{F}_t^L \subseteq \mathfrak{F}_t^X$ and $X_t = \{ all \bigcirc to left of L_t, all \bullet to right \}.$

Then,

$$\begin{split} \mathsf{E}\left[h(\Psi_t) \mid \mathfrak{F}_t^L\right] &= \mathsf{E}\left[\mathsf{E}\left[h(\Psi_t) \mid \mathfrak{F}_t^X\right] \mid \mathfrak{F}_t^L\right] \\ &= \mathsf{E}\left[\int h(\psi) \mathsf{\Pi}_{X_t}(d\psi) \mid \mathfrak{F}_t^L\right] \\ &= \int h(\psi) \mathsf{\Pi}_{L_t}(d\psi) \end{split}$$

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What Structure Is Inherited?

Space/type reversal of particle system has same distribution:



Therefore, $L_t = ^d - L_t$.

What Structure Is Inherited?

Conditional distribution $\Psi_{t+\cdot} \mid \mathfrak{F}_t^L$ is same as Ψ_{\cdot} shifted spatially.



Therefore, L_t has stationary, independent increments.

This Technique Generalizes

What is L?

The process L has:

- continuous paths;
- spatially symmetric distribution;
- stationary, independent increments.

So, L is a continuous, driftless Lévy process. Sounds like a Brownian motion! Let the particle motion (\hat{Z}, \hat{P}^z) be any continuous, Borel process with dual (Z, P^z) :

$$\int f P_t g dm = \int g \hat{P}_t f dm$$

for m having support on an interval of \mathbb{R} . Lay down initial particles using marginal location measure m.

Under some regularity conditions,

- ancestral trajectory is a copy of Z;
- initial interface persists as a point *L_t*;
- process *L* is continuous;
- $\mathfrak{L}(\Psi_t \mid \mathfrak{F}_t^L) = \Pi_{L_t};$
- L is a time-homogeneous Markov process.

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Markov Property

Since $L_t \equiv h(\Psi_t)$ and Ψ_t is a time-homogeneous Markov process with semigroup \mathbb{P}_t ,

$$\begin{split} \mathsf{E}\left[f(L_t) \mid \mathfrak{F}_s^L\right] &= \mathsf{E}\left[\mathsf{E}\left[f \circ h(\Psi_t) \mid \mathfrak{F}_s^\Psi\right] \mid \mathfrak{F}_s^L\right] \\ &= \mathsf{E}\left[\mathbb{P}_{t-s}f \circ h(\Psi_s) \mid \mathfrak{F}_s^L\right] \\ &= \int \mathbb{P}_{t-s}f \circ h(\psi) \mathsf{\Pi}_{L_s}(d\psi) \end{split}$$

so L_t is time-homogeneous with semigroup,

$$\mathsf{P}_t^L f(x) = \int \mathbb{P}_t f \circ h(\psi) \mathsf{\Pi}_x(d\psi)$$

A Duality Result

Theorem 4. $\mathsf{P}^l(L_t > y) = \mathsf{P}^y(Z_t < l)$

Proof. Follows immediately by a duality result in DEFKZ, 2000. $\hfill \Box$

Corollary 5. For \hat{Z} standard Brownian, the interface is standard Brownian.

Corollary 6. If Z has generator A and L has generator A_L , then

$$\int f' A_L g + \int g' A f = 0$$

General Diffusion

A general diffusion with generator

$$\hat{A}f(x) = \frac{a(x)}{2}f''(x) + b(x)f'(x)$$

is "usually" self-dual with respect to its speed measure

$$m(dx) = m_0 e^{2\int_0^x \frac{b(y)}{a(y)} dy} a^{-1}(x) dx$$

So, by Corollary 6,

$$A_L g(x) = \frac{a(x)}{2} g''(x) + \left(\frac{a'(x)}{2} - b(x)\right) g'(x)$$

In particular, if $a(x) \equiv a > 0$, general b(x) then A_L is diffusion with same diffusion rate and opposite drift.

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Interesting Examples



- $m(dx) = x^{-(2c+1)}dx$
- Generators:

$$\hat{A}f(x) = \frac{1}{2}x^2 f''(x) + \left(\frac{1}{2} - c\right)xf'(x)$$
$$\hat{A}_Lg(x) = \frac{1}{2}x^2g''(x) + \left(\frac{1}{2} + c\right)xg'(x)$$

• Processes:

$$Z_t = Z_0 e^{B_t - ct}$$
$$L_t = L_0 e^{B_t + ct}$$

• SPDEs:

$$Z_{t} = Z_{0} + \int_{0}^{t} Z_{s} dB_{s} + \left(\frac{1}{2} - c\right) \int_{0}^{t} Z_{s} ds$$
$$L_{t} = L_{0} + \int_{0}^{t} L_{s} dB_{s} + \left(\frac{1}{2} + c\right) \int_{0}^{t} L_{s} ds$$

Ornstein-Uhlenbeck

- $m(dx) = e^{-x^2/2}dx$
- Generators:

$$\hat{A}f(x) = \frac{1}{2}f''(x) - \frac{1}{2}xf'(x)$$
$$\hat{A}_{L}g(x) = \frac{1}{2}g''(x) + \frac{1}{2}xg'(x)$$

• Processes:

$$Z_{t} = e^{-t/2} \left(Z_{0} + B \left(e^{t} - 1 \right) \right)$$
$$L_{t} = e^{t/2} \left(L_{0} + B \left(1 - e^{-t} \right) \right)$$

• SPDEs:

$$Z_{t} = Z_{0} + B_{t} - \int_{0}^{t} \frac{1}{2} Z_{s} ds$$
$$L_{t} = L_{0} + B_{t} + \int_{0}^{t} \frac{1}{2} L_{s} ds$$

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Another Interesting Connection

The deterministic solution of

$$\begin{split} \dot{u} &= Au\\ u(0,x) &= \mathbf{1}_{x \geq L_0}\\ \text{is } u(x,t) &= \mathsf{P}^x(Z_t \geq 0) = \mathsf{P}^0(L_t \leq x). \end{split}$$

That is $u_t(\cdot)$ is the distribution function of L_t .